

UC Davis

UC Davis Previously Published Works

Title

Dynamics of epidemic diseases on a growing adaptive network.

Permalink

<https://escholarship.org/uc/item/4dh7b7d2>

Journal

Scientific reports, 7(1)

ISSN

2045-2322

Authors

Demirel, Güven
Barter, Edmund
Gross, Thilo

Publication Date

2017-02-01

DOI

10.1038/srep42352

Peer reviewed

SCIENTIFIC REPORTS

OPEN

Dynamics of epidemic diseases on a growing adaptive network

Güven Demirel¹, Edmund Barter² & Thilo Gross²

Received: 12 October 2016

Accepted: 08 January 2017

Published: 10 February 2017

The study of epidemics on static networks has revealed important effects on disease prevalence of network topological features such as the variance of the degree distribution, i.e. the distribution of the number of neighbors of nodes, and the maximum degree. Here, we analyze an adaptive network where the degree distribution is not independent of epidemics but is shaped through disease-induced dynamics and mortality in a complex interplay. We study the dynamics of a network that grows according to a preferential attachment rule, while nodes are simultaneously removed from the network due to disease-induced mortality. We investigate the prevalence of the disease using individual-based simulations and a heterogeneous node approximation. Our results suggest that in this system in the thermodynamic limit no epidemic thresholds exist, while the interplay between network growth and epidemic spreading leads to exponential networks for any finite rate of infectiousness when the disease persists.

Throughout human history epidemic diseases have been a constant threat. The plague of Athens killed 25–35 percent of the city's population as early as 430 BC¹. A bubonic plague epidemic killed between 75000 and 100000 inhabitants of London between April 1665 and January 1666, when its population was about 460000¹. Smallpox became a major threat to Europe throughout the eighteenth century, where growing cities like London were especially vulnerable to the disease on account of continuously providing immigrants that are susceptible to the virus¹. After a brief respite during the 20th century epidemics are on the rise again. For instance, around 655000 people died of malaria in 2010² and more than 100000 confirmed cases of Zika infection have occurred in the Americas in the past year³. Combating such epidemic diseases efficiently in the mega-cities of the future is likely to require a heightened understanding of the dynamics of the diseases and the social networks on which they spread.

In the past two decades, epidemic spreading has been extensively studied on different complex networks to understand the influence of the social contact network structure on the disease prevalence^{4–7}. Two commonly studied models of infectious diseases are Susceptible-Infected-Susceptible (SIS) and Susceptible-Infected-Removed (SIR) models.

For SIS diseases on complex networks, the crucial determinant of epidemic spreading is the average maximum degree or the degree cutoff of the contact network where the hub node with the maximum degree and its neighbors act as an active incessant source for the endemic state^{8,9}. Therefore, the epidemic threshold vanishes in the thermodynamic limit in networks where the maximum degree increases with the system size, which is true for most random networks unless there is a natural restriction on degree, for instance, due to cognitive limitations⁷.

For SIR diseases on complex networks, differently from SIS diseases, the fate of the disease is determined by the *degree distribution*, i.e. the probability distribution of finding an individual with a given number of contacts^{8,10,11}. If the variance of this distribution is finite, then there is generally a threshold of infectiousness, which the disease has to exceed to reach a finite fraction of the population. By contrast, if the variance of the degree distribution is infinite, such as in certain scale-free networks, then the epidemic threshold vanishes and the SIR disease percolates on the network for all non-zero infectiousness levels via hierarchical spreading from hub nodes to lower degree nodes¹². Thus, scale-free networks allow unlikely diseases with low infectiousness to spread and become endemic.

Today, the structure of social networks is recognized as a key factor with direct implications for epidemic dynamics and potential counter measures^{13–17}. This insight has motivated the integration of real world network data into epidemic models^{18–23}. Furthermore, theoretical models have been extended by including several properties of real-world networks such as degree constraints²⁴, degree correlations²⁵, clustering^{26–28}, information filtering²⁹, social hierarchy³⁰, and nonuniform transmission probabilities³¹.

¹Management Science & Entrepreneurship, Essex Business School, University of Essex, Southend-on-Sea, UK.

²Department of Engineering Mathematics, University of Bristol, Bristol, UK. Correspondence and requests for materials should be addressed to E.B. (email: edmund.barter@bristol.ac.uk)

A relatively recent addition is to consider also the feedback of the epidemic on the social network structure, which can be indirect, e.g. by triggering behavioral changes of agents^{6,32}, or direct by removing agents due to hospitalization, quarantine or death.

Modeling the network response to an ongoing epidemic leads to an adaptive network, a system with interplay between the dynamics of the network and the dynamics on the network takes place^{33,34}. Epidemics on adaptive networks can exhibit complex emergent dynamics^{32,35–38} (e.g. sustained oscillations, bistability, and hysteresis) and emergent topological properties^{32,35,39,40} (e.g. heterogeneous degree distributions and assortative degree correlations). Furthermore, the study of the social response to epidemics is interesting from an applied point of view because it could enable enhanced vaccine control⁴¹ and effective quarantine strategies⁴².

In comparison to social responses to epidemics^{32,35–39,43–48}, network growth and direct topological feedback via the removal of nodes have received less attention. Previous works considered the case where network growth and death processes are balanced and thus the population stays in equilibrium and fluctuates around a fixed system size^{49,50}. The case of continuous growth was studied in Ref. 51, where new nodes attach preferentially to high degree non-infected nodes. It was observed that a transition from a scale-free topology to an exponential one takes place as the infectiousness decreases. Reference 52 study the effects of network growth and demographics on the dynamics of an SIS disease simultaneously spreading on the network, showing that the epidemic threshold vanishes in the thermodynamic limit. In another study, network growth and node removals have also been incorporated in a single model⁵³. However, the authors focused on epidemic oscillations and did not consider topological effects in detail. Another related work focused on the interplay between network growth and dynamical behavior in the context of evolutionary game theory⁵⁴, where new players preferentially attach to those receiving higher payoffs.

Here, we consider the growth of a network by preferential attachment from which nodes are simultaneously removed due to an SIR epidemic. The appeal of this model lies in its paradoxical nature, in absence of the disease, preferential attachment leads to the formation of scale-free topologies in which the epidemic threshold vanishes, such that the disease can invade. However, an established SIR disease will quickly infect and remove nodes of high degree such that the variance of the degree distribution is decreased and epidemic thresholds reappear, potentially leading to the extinction of the disease.

Now, consider the following line of reasoning: Observing a scale-free topology implies that the epidemic is extinct. But an extinct epidemic implies scale-free structure and hence vanishing epidemic threshold precluding extinction. Logically, the only possible solution is that the epidemic persists (unconditionally) in a network that is not scale-free. In other words, one would expect that the coevolution of epidemic state and network structure should lead to a vanishing epidemic threshold in an exponential network.

The argument presented above is admittedly hand-wavy. One objection against this line of reasoning that comes to mind immediately is that the paradox can also be resolved temporally, such that the epidemic goes extinct while the network is exponential. A disease-free scale-free network can then develop later since the disease cannot be reintroduced since no infected agents are left. However, this temporal resolution is only feasible in finite networks. In the thermodynamic limit it can be easily shown that a finite number of infected survive even below the epidemic threshold, which precludes complete extinction.

In the remainder of this paper we present a detailed dynamical analysis of the epidemic model using a heterogeneous node approximation along with detailed numerical computations. We show that disease-induced mortality reduces the variance of the degree distribution to a finite value, but not sufficiently far to cause the extinction of the epidemic. Thus, a balance is reached where the finite variance of the degree distribution is matched to the infectiousness of the pathogen. Therefore, the epidemic threshold vanishes in the thermodynamic limit, confirming the hand-waving argumentation above. We also identify a parameter region, where, after all, a temporal resolution of the paradox is observed.

Methods

In this section, we first introduce the model and then develop our theoretical approach that describes the evolution of the network dynamics based on coarse-graining approximations.

Model. We study the spreading of a susceptible-infected-removed (SIR) disease⁵⁵ on an evolving network. In this network, a given node is either susceptible (state S) or already infected with the disease (state I). We start with a fully connected network of m_0 nodes and consider three dynamical processes: a) the arrival of nodes, b) disease transmission, and c) the removal of nodes.

In the following we measure all rates per capita, including the arrival rate. This implies that larger populations have a proportionally larger influx of individuals, which appears plausible e.g. for growing cities, where the attractiveness of the city increases with size. It is analogous to the use of per capita birth rates in models of population dynamics. We note that this assumption is necessary to keep the model well-defined in the thermodynamic limit.

New nodes arrive in the population at a constant per capita rate q and are already infected with the disease with probability w . Arriving nodes immediately establish links with m of the nodes selected according to the preferential attachment rule⁵⁶: A new node establishes a link with a particular node i of degree k_i with a probability proportional to $k_i/\sum_j k_j$. Therefore, an incoming node's links will attach to a node of degree k with probability $kp_k/\langle k \rangle$, where p_k denotes the degree distribution (the probability that a randomly picked node has degree k) and $\langle k \rangle = \sum k p_k$ is the mean degree.

Disease transmission occurs at rate p on every link connecting a susceptible and an infected node. Therefore, nodes with higher degree are proportionally more likely to catch and spread the disease.

Removal of infected nodes takes place at rate r . Because we describe a fatal disease from which recovery is not possible, removed nodes and their links are entirely deleted and do not re-appear at a later stage. The removal of links along with the nodes depicts the rapid removal of corpses in the human population. The removal

mechanism that is used here can also be considered as an approximate description for the hospitalization of infected individuals, which effectively removes the links. Furthermore, diseases in which dead hosts continue transmitting the disease can be captured in the same framework by a reduced removal rate r , taking into account the finite time between the death and the actual removal of the body.

Unless mentioned otherwise, the following set of parameter values is used throughout the paper: $m_0 = 6$, $m = 5$, $q = 0.01$. In agent-based simulations the network is simulated until N reaches 10^7 or the time reaches 10^4 .

Analytical treatment. The dynamics on and of complex networks can be captured by a set of coupled ordinary differential equations, in so-called coarse-graining or moment-closure approximations^{7,39,57–64}. In the next section we develop a heterogeneous node approximation, also called heterogeneous mean-field or degree-based mean-field approximation^{7,11,65,66}, where the network evolution is captured in a set of equations for the node densities in different degree-classes.

Since the heterogeneous node approximation is in the form of a high-dimensional system of ordinary differential equations, we follow two directions for reducing the dimensionality of analysis. First we apply the mathematical triple jump approach⁶⁷ to transform the infinite dimensional ordinary differential equation system in the thermodynamic limit to a two-dimensional partial differential equation system. Second we develop an alternative approach that reduces the heterogeneous node approximation to a low-dimensional ordinary differential equation system assuming random graph properties, which is later in the paper shown to be capable of estimating the network dynamics to high accuracy when away from the epidemic threshold.

Heterogeneous approximation. The heterogeneous node approximation consists of a set of ordinary differential equations for the densities $[A_k]$, the abundance of nodes in the class A_k , which is the set of nodes of state $A \in \{S, I\}$ and degree k , normalized by the total number of nodes N . The total density of S -nodes is denoted as $[S]$, $[S] = \sum_k [S_k]$. The density $[I]$ is defined analogously such that $[S] + [I] = 1$.

Before deriving the full equations for the specific system under investigation, we first illustrate the general structure of the heterogeneous moment expansion

$$\frac{d}{dt}[A_k] \equiv \lim_{\Delta t \rightarrow 0} \frac{[A_k(t + \Delta t)] - [A_k(t)]}{\Delta t}.$$

Here, two types of terms contribute to $d[A_k]/dt$: a) changes in the abundance of nodes in the class A_k , and b) changes in the normalization factor N .

For illustration, let us consider the process that removes infected individuals at a per capita rate r . For the moment we assume that these removals do not change N , as we treat the change in N separately below. When an infected node of degree k is removed, the abundance of nodes of type I_k , $N_{(I,k)}$, decreases by 1 so that the density $[I_k]$ is reduced by $1/N$. Considering that $(\Delta t)rN_{(I,k)}(t)$ such removal events take place within a time step Δt , we obtain the rate of change for the density $[I_k]$ due to process a):

$$\frac{d}{dt}[I_k] = -r \frac{N_{(I,k)}(t)}{N(t)} = -r[I_k].$$

We now express the effect of the modification of the normalization factor N due to such removal events. When an infected node of arbitrary degree k' is removed, the densities of all degree classes A_k are affected due to the modified normalization factor. In total, $(\Delta t)rN_I$ removal events take place within a time step Δt resulting in $[A_k(t + \Delta t)] = [A_k(t)]N(t)/(N(t) - 1)$, when isolated from the other changes of type a). Therefore, the rate of change of type b) due to removal events is

$$\begin{aligned} \frac{d}{dt}[A_k] &= \lim_{\Delta t \rightarrow 0} rN_I(t) \left([A_k(t)] \frac{N(t)}{N(t) - 1} - [A_k(t)] \right) \\ &= \lim_{\Delta t \rightarrow 0} r[I(t)][A_k(t)] \frac{N(t)}{N(t) - 1}. \end{aligned}$$

Noting that every node is updated on average once in a unit time, i.e. $\Delta t = 1/N$, and taking the thermodynamic limit, we obtain the renormalization rate

$$\frac{d}{dt}[A_k] = r[I][A_k].$$

Writing the complete set of changes caused by infection, node arrival, and node removal processes, we derive the moment expansion for the densities $[S_k]$ and $[I_k]$

$$\begin{aligned}
\frac{d}{dt}[S_k] &= q((1-w)\delta_{k,m} + \frac{m}{\langle k \rangle}(-k[S_k] + (k-1)[S_{k-1}] - [S_k]) - p\sum_{k'}[S_k I_{k'}] \\
&\quad + r(\sum_{k'}[S_{k+1} I_{k'}] - \sum_{k'}[S_k I_{k'}] + [I][S_k]), \\
\frac{d}{dt}[I_k] &= q(w\delta_{k,m} + \frac{m}{\langle k \rangle}(-k[I_k] + (k-1)[I_{k-1}] - [I_k]) + p\sum_{k'}[S_k I_{k'}] \\
&\quad + r(\sum_{k'}[I_{k+1} I_{k'}] - \sum_{k'}[I_k I_{k'}] + [I][I_k] - [I_k]), \quad 0 < k < k_M
\end{aligned} \tag{1}$$

here, the quantity $[A_k B_{k'}]$ denotes the density of links between nodes of type A_k and $B_{k'}$, where $A, B \in \{S, I\}$, and k and k' are the respective degrees.

For understanding the equation governing the evolution of the density $[S_k]$, consider that new nodes with degree m arrive at rate q and have state S with probability $1-w$. Thus, the density $[S_m]$ increases at rate $q(1-w)$. A newly arriving node builds a link to a node in the S_k class with probability $k[S_k]/\langle k \rangle$ and causes it to pass into the S_{k+1} class. Because m such links are established by each newly arriving node, the density $[S_k]$ decreases by $qmk[S_k]/\langle k \rangle$. Similarly, nodes in the S_{k-1} class pass into the S_k class at rate $qm(k-1)[S_{k-1}]/\langle k \rangle$. Additionally, as explained above, we need to renormalize the density $[S_k]$ when a node arrives. This corresponds to a loss of the $[S_k]$ density of $q[S_k]$.

At rate p , nodes within the S_k class become infected through their links with infected nodes of arbitrary degree k' , causing them to pass into the I_k class. The total density of such links is $\sum_{k'}[S_k I_{k'}]$.

Finally, nodes within the S_k class pass into the S_{k-1} class due to the removal of their infected neighbors of arbitrary degree k' . Given the density of such links, the density of $[S_k]$ decreases by $r\sum_{k'}[S_k I_{k'}]$. Similarly, nodes in the S_{k+1} class pass into the S_k class corresponding to a gain of $r\sum_{k'}[S_{k+1} I_{k'}]$. As infected nodes are removed, the density of all degree classes increases due to the renormalization leading to a gain of $r[I][S_k]$ for the density $[S_k]$. The rate equation for the density $[I_k]$ is constructed analogously, with the addition of a term for the removal of infected nodes with degree k at the rate $r, r[I_k]$.

As the equations for node densities $[A_k]$ depend on link densities $[A_k B_{k'}]$, Eq. (1) is not closed. In order to close the system, the moment expansion should be truncated by the moment-closure approximation, in which the densities of larger subgraphs are estimated in terms of the densities of smaller ones. Here, we use the heterogeneous node approximation to close the system at the node level

$$[A_k B_{k'}] \approx \frac{kk'[A_k][B_{k'}]}{\langle k \rangle}. \tag{2}$$

We have assumed that the nodes with the same degree can be considered identical and state and degree correlations between neighboring nodes are negligible. The mixing assumption generally requires a mixing or annealing process that makes it possible to replace the adjacency matrix structure with the degree distribution⁸, which is provided here by the constant removal and addition of nodes and links.

Using the node approximation of Eq. (2), we reach

$$\begin{aligned}
\frac{d}{dt}[S_k] &= q((1-w)\delta_{k,m} + \frac{m}{\langle k \rangle}(-k[S_k] + (k-1)[S_{k-1}] - [S_k]) - pz_I I k[S_k] \\
&\quad + r(z_I I((k+1)[S_{k+1}] - k[S_k]) + [I][S_k]), \\
\frac{d}{dt}[I_k] &= q(w\delta_{k,m} + \frac{m}{\langle k \rangle}(-k[I_k] + (k-1)[I_{k-1}] - [I_k]) + pz_I I k[S_k] \\
&\quad + r(z_I I((k+1)[I_{k+1}] - k[I_k]) + [I][I_k] - [I_k]), \quad 0 < k < k_M
\end{aligned} \tag{3}$$

where $z_I = \langle k_I \rangle / \langle k \rangle$ and $\langle k_I \rangle$ is the mean degree of infected nodes.

In the following we refer to Eq. (3) as the heterogeneous approximation. The main drawback of such heterogeneous approximations is the high dimensionality of the system of equations, which complicates the analytical solution and thus typically necessitates extensive numerical studies except for the analysis of special conditions. Furthermore, since it is not possible to numerically integrate an infinite dimensional system of differential equations, we need to introduce a degree cut-off k_M by assuming $\sum_{k=k_M+1}^{\infty} p_k \ll \sum_{k=0}^{k_M} p_k$. The higher the degree cut-off, k_M , the more precise the heterogeneous approximation becomes.

Mathematical triple jump approach. The mathematical triple jump approach of ref. 67 consists of three steps. First, a high-dimensional ordinary differential equation system is developed to capture the dynamics under the types of heterogeneity that are identified to be of utmost importance. Second, the obtained system of ordinary differential equations is transformed to a low-dimensional partial differential equation system in the thermodynamic limit using moment generating functions. Finally, the partial differential equation system is analyzed using the tools of dynamical systems theory.

The first step has already been completed to find the heterogeneous approximation in Eq. (3). The second step is done below in this section, while the last step is carried out in the next section.

We first introduce the quantities $Q(t, x) = \sum_k [S_k(t)] x^k$ and $R(t, x) = \sum_k [I_k(t)] x^k$, which are the generating functions of the degree distributions of susceptible and infected populations. The time derivatives of $Q(t, x)$ and $R(t, x)$ are given by

$$Q_I(t, x) \equiv \sum_k \frac{d[S_k(t)]}{dt} x^k R_I(t, x) \equiv \sum_k \frac{d[I_k(t)]}{dt} x^k$$

The partial derivatives with respect to x are defined analogously by

$$Q_x(t, x) \equiv \sum_k [S_k(t)] k x^{k-1} R_x(t, x) \equiv \sum_k [I_k(t)] k x^{k-1}$$

The functions $Q(t, x)$ and $R(t, x)$ are particularly useful because they are related to the moments as given below:

$$\gamma \equiv R(t, 1) = [I] = 1 - [S], \quad \alpha \equiv Q_x(t, 1) = \langle k_S \rangle [S], \quad \beta \equiv R_x(t, 1) = \langle k_I \rangle [I].$$

Using these quantities, we obtain the partial differential equations

$$Q_t(t, x) = q(1 - w)x^m + (r\gamma - q)Q(t, x) + \frac{1}{\alpha + \beta} [qmx^2 - (qm + (p + r)\beta)x + r\beta] Q_x(t, x) \quad (4)$$

$$R_t(t, x) = qwx^m - [q + r(1 - \gamma)]R(t, x) + \frac{1}{\alpha + \beta} (1 - x)(r\beta - qmx)R_x(t, x) + \frac{p\beta x}{\alpha + \beta} Q_x(t, x). \quad (5)$$

Homogeneous approximation. As a second alternative approach we develop a low-dimensional approximation by summing over the degree classes in Eq. (3). We consider the susceptible proportion of the population $[S]$, the mean degree $\langle k \rangle$, and the mean degree of susceptibles $\langle k_S \rangle$ which evolve according to

$$\begin{aligned} \frac{d}{dt}[S] &= \sum_k \frac{d}{dt}[S_k], \\ \frac{d}{dt}\langle k \rangle &= \sum_k k \frac{d}{dt}([S_k] + [I_k]), \\ \frac{d}{dt}\langle k_S \rangle [S] &= \sum_k \frac{d}{dt}k[S_k]. \end{aligned} \quad (6)$$

Using Eqs (3) and (6), we obtain

$$\begin{aligned} \frac{d}{dt}[S] &= q(1 - w - [S]) - p \frac{\langle k_S \rangle \langle k_I \rangle}{\langle k \rangle} [S][I] + r[S][I], \\ \frac{d}{dt}\langle k \rangle &= q(2m - \langle k \rangle) + r(2\langle k_S \rangle [S] - \langle k \rangle(1 + [S])), \\ \frac{d}{dt}\langle k_S \rangle &= q \left(\frac{(1 - w)(m - \langle k_S \rangle)}{[S]} + m \frac{\langle k_S \rangle}{\langle k \rangle} \right) - p \frac{\langle k_I \rangle [I]}{\langle k \rangle} (\langle k_S^2 \rangle - \langle k_S \rangle^2) \\ &\quad - r \frac{\langle k_S \rangle \langle k_I \rangle}{\langle k \rangle} [I]. \end{aligned} \quad (7)$$

In the following, we refer to Eq. (7) as the coarse-grained heterogeneous approximation, which does not involve any further approximations other than those discussed above.

Because we have not derived an equation for the second moment of the susceptible degree distribution $\langle k_S^2 \rangle$, Eq. (7) does not constitute a closed dynamical system. We address this problem by replacing $\langle k_S^2 \rangle$ by $\langle k_S \rangle^2 + \langle k_S \rangle$ in an additional approximation. We note that this approximation is valid exactly when the network has a Poisson degree distribution. It can therefore be thought of as a ‘random-graph approximation’. This approximation will certainly fail in the case of scale-free networks because of the degree distribution’s diverging variance, i.e. $\langle k_S^2 \rangle \rightarrow \infty$, in which case we will resort to the heterogeneous approximation and its PDE description in the thermodynamic limit for the analysis of the model. However, as will become apparent below, the system obtained by the random-graph approximation still performs well for distributions with large finite variance.

Using the random-graph approximation we obtain

$$\begin{aligned} \frac{d}{dt}[S] &= q(1 - w - [S]) - p \frac{\langle k_S \rangle \langle k_I \rangle}{\langle k \rangle} [S][I] + r[S][I], \\ \frac{d}{dt}\langle k \rangle &= q(2m - \langle k \rangle) + r(2\langle k_S \rangle [S] - \langle k \rangle(1 + [S])), \\ \frac{d}{dt}\langle k_S \rangle &= q \left(\frac{(1 - w)(m - \langle k_S \rangle)}{[S]} + m \frac{\langle k_S \rangle}{\langle k \rangle} \right) - p \frac{\langle k_S \rangle \langle k_I \rangle [I]}{\langle k \rangle} - r \frac{\langle k_S \rangle \langle k_I \rangle}{\langle k \rangle} [I], \end{aligned} \quad (8)$$

where $[I]$ and $\langle k_I \rangle$ are given by the equations $[S] + [I] = 1$ and $\langle k_S \rangle [S] + \langle k_I \rangle [I] = \langle k \rangle$, such that the system constitutes a closed model. In the following we refer to this model as the homogeneous approximation.

Results

In the previous section, we first present the analysis of heterogeneous and homogeneous approximations and confirm the results by comparison with individual-based simulations of the network. Later we present a detailed analysis of the epidemic threshold. We finally discuss the emergence of dynamics that involve epidemic cycles.

General properties of the network and disease prevalence. Before we launch into a detailed discussion of the model, let us consider the limiting case of network evolution in the absence of the epidemic. In this case the model is identical to the Barabási-Albert model of network growth⁵⁶, which is known to lead to scale-free topologies, where the degree distribution follows a power law $p_k \propto k^{-\gamma}$ with exponent $\gamma = 3$ and thus the degree variance σ^2 diverges in the disease-free state. Because the density of infected vanishes in the absence of the epidemic, it is also evident that the degree distribution must be scale-free independent of the parameters p and r .

In the present model the emergence of scale-free topologies is thus expected in the limit where the disease goes extinct or remains limited to a finite number of infected nodes $\ll N$. When the epidemic is present, high degree nodes are disproportionately likely to become infected and subsequently removed, which can be expected to prevent the formation of scale-free topologies.

We confirm this intuition by plotting degree distributions for various parameter sets in Fig. 1. We show a comparison of the heterogeneous approximation with individual-based simulations. The figure shows a good agreement between the modeling approaches and confirms basic intuition. When all arriving nodes are susceptible ($w = 0$), a scale-free degree distribution with the expected exponent $\gamma = 3$ is formed for $p = 0$. At finite infectiousness p , the topology changes from scale-free to exponential. The same behavior is observed at higher rates of infected arrivals, $0 < w < 1$.

When all arriving nodes are already infected ($w = 1$), the distribution has a bimodal form for $p = 0$ with high degree contribution coming from the initial susceptibles which never get infected. At positive infectiousness p , these individuals eventually die and the mode at high degrees disappears.

In order to quantify the topological transition from the scale-free to the exponential degree distribution, we plot the variance σ^2 of the degree distribution as a function of infectiousness p and fraction of infected arrivals w in Fig. 2. As either parameter increases, the disease prevalence in the steady-state, $[I]^*$, increases and removal occurs at a high rate. As a result the degree distribution becomes narrower and the degree variance decreases. It is apparent that very high values of the variance are only found for low infectiousness p , whereas higher infectiousness quickly leads to narrow distributions.

Above we computed the variance σ^2 of the degree distribution of networks. One concern in any computation of this kind is finite-size effects. In the heterogeneous approximation these effects appear directly in form of the maximal degree that is considered in the approximation. In agent-based simulation a similar cut-off exists as the maximal degree in a network of finite size is bounded by the number of nodes. Hence all moments of the degree distribution, including the variance, must be finite regardless of the shape of the degree distribution. However, if the finite networks are drawn from an ensemble that becomes scale-free in the thermodynamic limit the variance σ^2 is often found to increase logarithmically with the imposed cut-off²⁴.

We now rule out that low values of the degree variance σ^2 , observed above, were due to finite-size effects by considering the variance σ^2 as a function of the degree cut-off k_c (Fig. 2 inset). For the case of $p = 0$, where we observed scale-free behavior, we find that the observed variance increases logarithmically as expected. Conversely, for the finite values of infectiousness p the observed σ^2 is insensitive to a sufficiently large cut-off. In summary these results show that fatal diseases should relatively quickly destroy scale-free structure of social networks at all but the smallest removal rate and/or infectiousness.

Let us now investigate the effect of the emergent network topology on the prevalence of the disease. Plots of the disease prevalence as a function of the infectiousness p and the fraction of infected arrivals w are shown in Figs 3 and 4. Figure 3 shows that the heterogeneous approximation (dashed lines) is in very good agreement with the agent-based model for the whole range of parameters p and w . The precision of the homogeneous approximation is also high for a large range of parameter values. As illustrated in Fig. 4, the absolute error in estimation of the disease prevalence of the approximation is maximal for intermediate values of infectiousness p , but still less than 0.05. The only qualitative discrepancy between the approximation and the agent-based model emerges at low infectiousness p for zero infected arrivals ($w = 0$). Here, the homogeneous model predicts the existence of an epidemic threshold, whereas in the agent-based simulation and the heterogeneous approximation the disease is found to persist even for very low levels of infectiousness p .

Summarizing the results shown so far, we can say that ongoing epidemic dynamics quickly leads to the formation of networks with finite variance. Generally, one would expect that such networks should exhibit a finite epidemic threshold. Nevertheless, the heterogeneous approximation and simulations indicate that the epidemic can persist in these networks for any finite positive value of the infectiousness.

Analysis of the epidemic threshold. Here, we investigate in greater detail the apparent absence of the threshold, which is implied by the general analysis above. Throughout the argument we will only consider the case where all arriving agents are susceptible, $w = 0$.

The epidemic threshold is commonly defined as the minimal value of the infectiousness, below which a randomly picked node is susceptible with probability 1. In any finite network this implies that below the epidemic threshold each individual node is susceptible. By contrast in the thermodynamic limit there can still be a finite number of infected nodes as long as the density of such nodes in the network is zero, i.e. $[I] = 0$. In the following

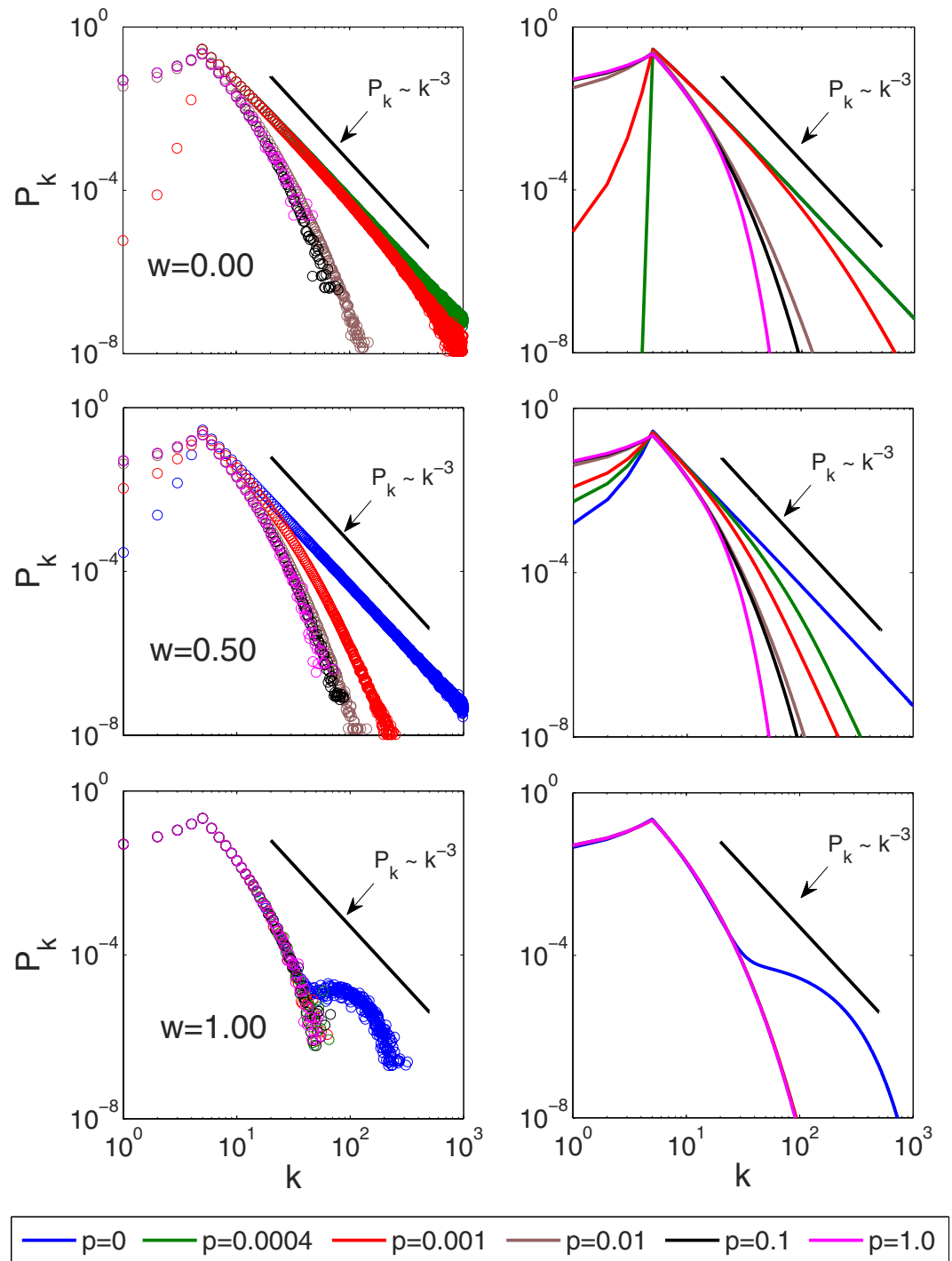


Figure 1. Degree distribution for varying infectiousness p and infected arrival fraction w . Left: individual-based simulations. Right: heterogeneous approximation. Scale-free degree distributions are observed when disease infectiousness vanishes ($p=0$). For increasing infectiousness the degree distribution quickly becomes exponential. Parameters: $r=q=0.01$, $m_0=6$, $m=5$. Shown are averages over 1000 simulation runs.

we denote the state below the epidemic threshold as the disease-free state, but recognize that there may be still a finite number of infected individuals.

We now calculate the epidemic threshold from Eq. (3) by looking at the stability of the disease-free state represented by $[I] = 0$ and hence $[I_k] = 0$ for all k . The disease-free state is stable if all eigenvalues of the corresponding Jacobian matrix have negative real parts. The epidemic threshold is characterized by a bifurcation point where the leading eigenvalue of the Jacobian is zero. The disease-free state becomes unstable while a stationary state with non-zero disease prevalence becomes stable beyond the epidemic threshold.

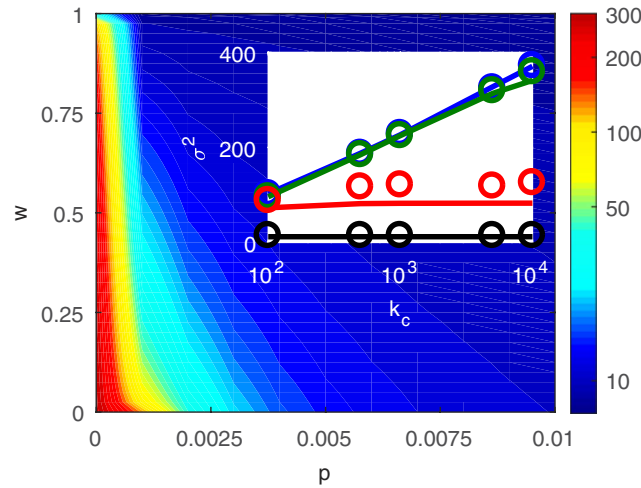


Figure 2. Variance σ^2 of the degree distribution. Degree variance σ^2 is plotted as a function of the infectiousness p and the infected arrival fraction w from agent-based simulations. The degree variance σ^2 decreases with increasing infectiousness p and fraction of infected arrivals w . Inset: Degree variance σ^2 as a function of the degree cut-off k_c . A transition from a cut-off dependent ($p = 0$, scale-free) to an independent ($p = 0.0004$, $p = 0.001$, and $p = 0.005$, exponential) regime is observed as the infectiousness p is increased. In network simulations (circles), nodes were restricted to at most k_c neighbors. In the heterogeneous approximation (solid lines) the cut-off k_c is directly imposed as k_M . Parameters: $r = q = 0.01$, $m_0 = 6$, $m = 5$, $k_c = 5 \times 10^3$.

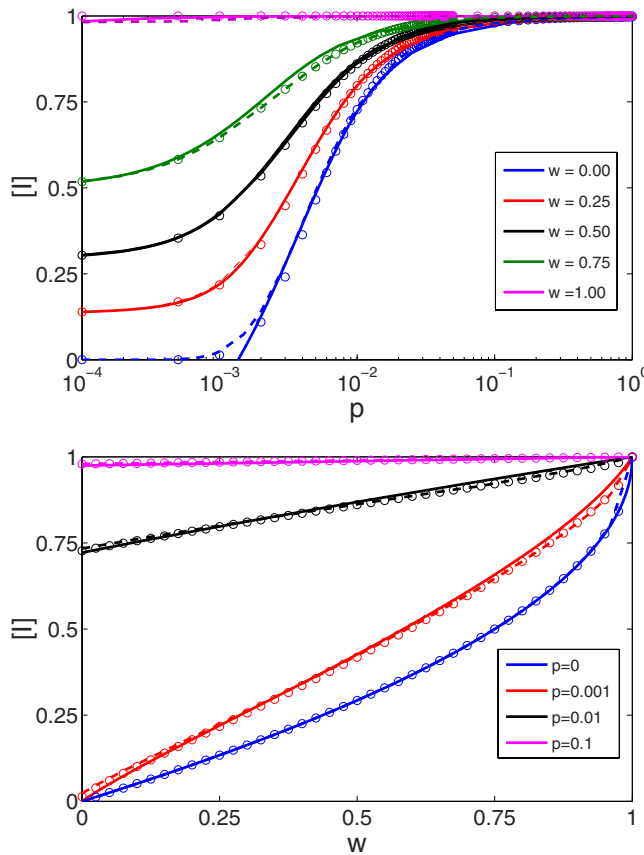


Figure 3. Dependence of the disease prevalence $[I]^*$ on infectiousness p (top) and fraction of infected arrivals w (bottom). The disease prevalence increases monotonically with infectiousness p and fraction of infected arrivals w . In agent-based simulations (circles), $[I]^*$ is calculated over the surviving runs among 10^3 total realizations. Homogeneous approximation (solid lines) is the analytical solution of Eq. (8). Heterogeneous approximation (dashed lines) is the stationary value of the numerical integration of Eq. (3). Parameters: $r = q = 0.01$, $m_0 = 6$, $m = 5$.

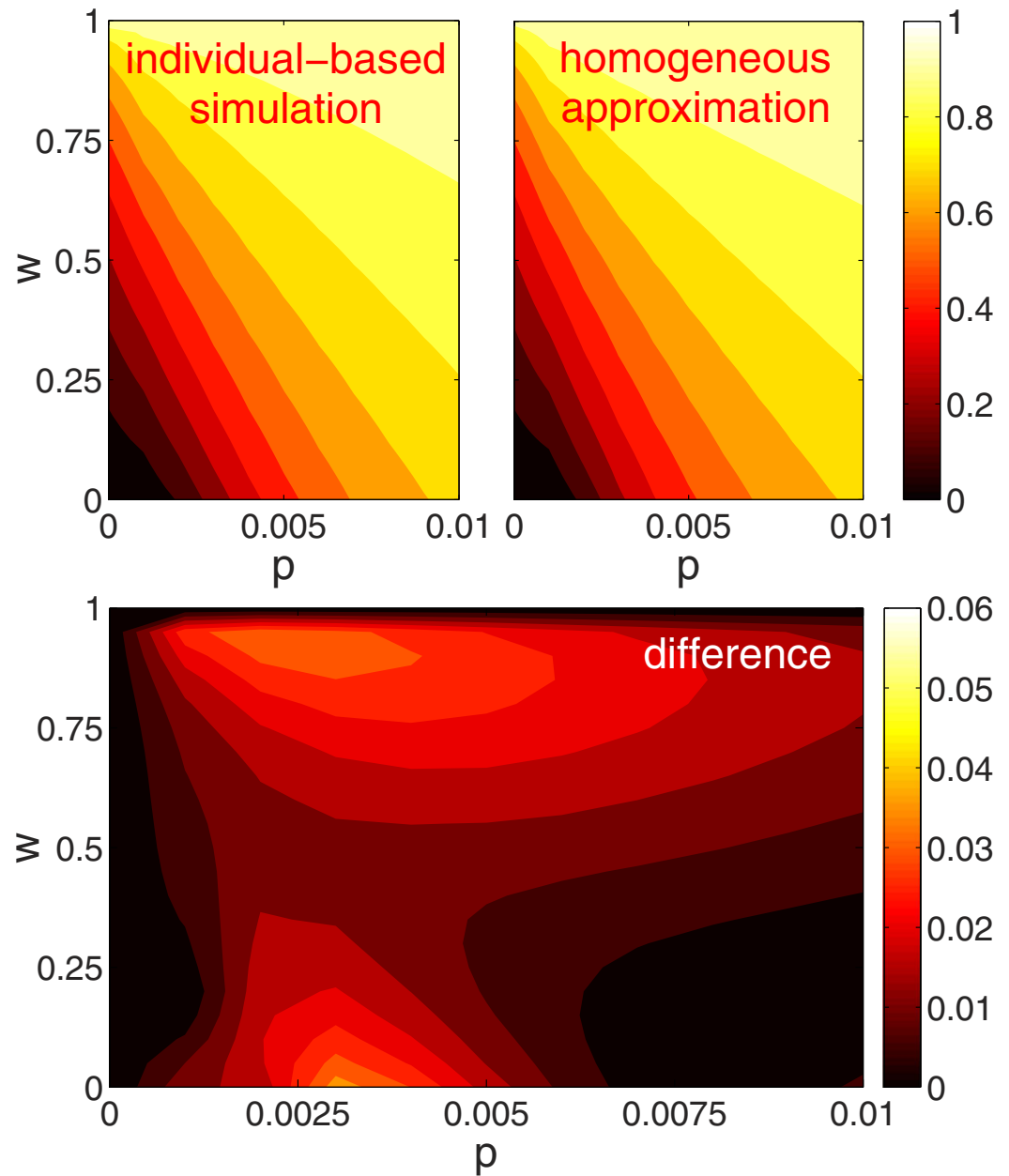


Figure 4. Performance of the homogeneous approximation in estimating the disease prevalence $[I]^*$ in the (p, w) parameter space. Top-left: individual-based simulations. Top-right: homogeneous approximation. Bottom: absolute difference between individual-based simulations and the homogeneous approximation. Homogeneous approximation performs well for a large range of infectiousness p and infected arrival fractions w . Parameters: $r = q = 0.01$, $m_0 = 6, m = 5$.

The Jacobian matrix is obtained by the linearization of Eq. (3) around the steady state. We will shortly refer to the Jacobian matrix evaluated at the disease-free steady state as the Jacobian and denote it $J = [J_{ij}]$, where $i, j = 0, 1, 2, \dots, k_M$. The Jacobian has a block matrix structure

$$J = \begin{pmatrix} \Theta_S & 0 \\ C & \Theta_I \end{pmatrix} \quad (9)$$

where $\Theta_S \equiv \partial_S dS/dt$, $\Theta_I \equiv \partial_I dI/dt$, and $C \equiv \partial_I dS/dt$, all evaluated at the disease-free state and with the definitions $S \equiv ([S_0], [S_1], [S_2], \dots, [S_{k_M}])$ and $I \equiv ([I_0], [I_1], [I_2], \dots, [I_{k_M}])$.

The epidemic threshold is characterized by zero eigenvalue of the Jacobian, which is given by the solution of $\det J = 0$. From Eq. (9), the determinant of the Jacobian can be obtained from the equality of $\det J = \det \Theta_S \det \Theta_I$.

Taking the partial derivatives of Eq. (3), we obtain

$$\Theta_S = u_S^\top v_S + A_S, \quad (10)$$

where $u_S = [u_S(k)]$, $v_S = [v_S(k)]$, and $A_S = [A_S(j, k)]$, for $j, k = 1, 2, \dots, k_M + 1$, with matrix entries

$$u_S(k) = \begin{cases} 0, & k = 1 \\ \frac{q}{4m}((k-1)[S_{k-1}] - (k-2)[S_{k-2}]), & 2 \leq k \leq k_M \\ -\frac{q}{4m}(k-2)[S_{k-2}], & k = k_M + 1 \end{cases}$$

$$v_S(k) = k - 1$$

$$A_S(j, k) = \begin{cases} -\frac{q(j+1)}{2}, & j = k \\ \frac{q(j-1)}{2}, & j = k + 1 \\ 0, & \text{otherwise} \end{cases}$$

From the matrix determinant lemma, $\det \Theta_S = (1 + v_S A_S^{-1} u_S^\top) \det A_S$.

$$(1 + v_S A_S^{-1} u_S^\top) = 0, \quad (11)$$

which can be computed exactly given that, for $Q_S \equiv \frac{q}{2} A_S^{-1}$ and $Q_S = [Q_S(j, k)]$, the entries of Q_S are given by the recursion equation

$$Q_S(j, k) = \begin{cases} 0, & k > j \\ -\frac{1}{k+1}, & k = j \text{ and } j < k_M + 1 \\ -\frac{1}{2}, & k = j = k_M + 1 \\ \frac{k-1}{k+1} Q_S(j, k+1), & k < j \end{cases} \quad (12)$$

We note that the matrix Q_S provides also the densities $[S_k]$ in the disease-free steady state for finite degree cut-off k_M . The vector of steady state susceptible densities, S , can be solved from

$$S = -2Q_S b, \quad (13)$$

where

$$b_i = \begin{cases} 1, & i = m \\ 0, & \text{otherwise} \end{cases} \quad (14)$$

We have now established an explicit solution for $\det \Theta_S = 0$ for any k_M , which is obtained by inserting Eqs (12) and (13) into Eq. (11).

We now follow the same procedure to solve $\det \Theta_I = 0$, which constitutes the other possible solution of $\det J = 0$. This solution is given by

$$(1 + v_I A_I^{-1} u_I^\top) = 0, \quad (15)$$

where

$$u_I(k) = \frac{p}{2m}(k-1)S_{k-1}$$

$$v_I(k) = k - 1$$

$$A_I(j, k) = \begin{cases} -r - \frac{q(j+1)}{2}, & j = k \\ \frac{q(j-1)}{2}, & j = k + 1 \\ 0, & \text{otherwise} \end{cases}$$

The matrix inverse $Q_I \equiv A_I^{-1}$ is solved from the recursive equation

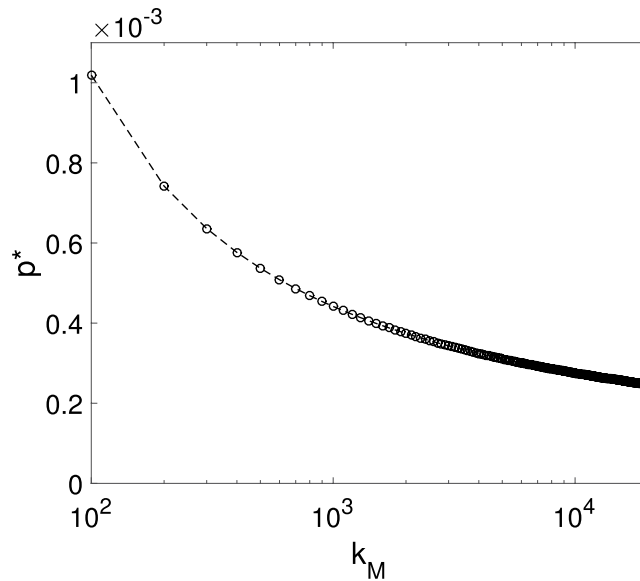


Figure 5. The epidemic threshold P^* calculated from Eq. (15) as a function of the degree cut-off k_M . The epidemic threshold vanishes in the thermodynamic limit. Parameters: $r = q = 0.01$, $w = 0$, $m_0 = 6$, $m = 5$.

$$Q_I(j, k) = \begin{cases} 0, & k > j \\ -\frac{1}{r + q(k+1)/2}, & k = j \\ \frac{q(k-1)}{2r + q(k+1)} Q_I(j, k+1), & k < j \end{cases} \quad (16)$$

The epidemic threshold needs to satisfy one of Eqs (11) or (15) for any finite degree cut-off k_M . Our investigations show that Eq. (11) is ruled out and the epidemic threshold is given by Eq. (15). Figure 5 shows the epidemic threshold calculated from Eq. (15) for increasing values of degree cut-off k_M , which approaches zero in the thermodynamic limit. This result is constraintuitive, as the epidemic threshold vanishes even in a system where the degree variance is finite.

For verification, we now take a complementary approach and calculate the epidemic threshold from Eqs (4) and (5). Setting time derivatives to 0 and $x = 1$ in both Eqs (4) or (5), we obtain

$$r\gamma^2 - (q + r)\gamma + \frac{p\beta\alpha}{\alpha + \beta} = 0.$$

for the case $w = 0$. Using the definitions in Eq. (4) we find

$$\left(r - p \frac{\langle k_I \rangle \langle k_S \rangle}{\langle k \rangle} \right) \gamma^2 + \left(p \frac{\langle k_I \rangle \langle k_S \rangle}{\langle k \rangle} - (q + r) \right) \gamma = 0.$$

As expected there are two solutions for the steady state fraction of nodes infected, γ . Factoring out the uninfected state $\gamma = 0$ gives the non-trivial root

$$\gamma = \frac{q + r - p \frac{\langle k_I \rangle \langle k_S \rangle}{\langle k \rangle}}{r - p \frac{\langle k_I \rangle \langle k_S \rangle}{\langle k \rangle}}. \quad (17)$$

For this root to give a value of $0 \leq \gamma \leq 1$ requires $(q + r) < p_c \frac{\langle k_I \rangle \langle k_S \rangle}{\langle k \rangle}$.

One can now ask if $(q + r) = p_c \frac{\langle k_I \rangle \langle k_S \rangle}{\langle k \rangle}$ constitutes a finite threshold below which the endemic state becomes unphysical. At such a threshold the number of infected nodes would become zero, leading to a double root at $\gamma = 0$. At this point $\langle k \rangle \approx \langle k_S \rangle$ and therefore

$$p_c \approx \frac{q + r}{\langle k_I \rangle}. \quad (18)$$

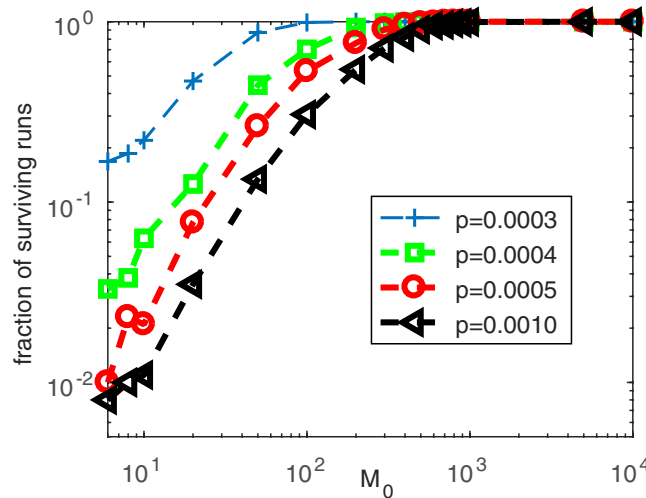


Figure 6. The fraction of surviving runs for low infectiousness p as a function of the initial network size M_0 . The fraction of surviving runs increases as the initial network size M_0 increases. In order to ensure the same initial average degree m , the Barabási-Albert growth model with $m = 5$ and $m_0 = 6$ is iterated until the network reaches size M_0 , then the full model simulation with disease dynamics starts. Then nodes are assigned states and the infection, removal, and network growth processes take place simultaneously. Parameters: $r = q = 0.01$, $w = 0$, $m_0 = 6$, $m = 5$.

We note that for large networks the probability of a node being infected scales as kp_k . Therefore in the case of the scale-free degree distribution, $p_k \sim k^{-3}$, which is bound to form at the threshold, if one existed, the mean degree of infected nodes, $\langle k_i \rangle$, is undefined. Hence in the limit of large networks $p_c \rightarrow 0$ and the threshold vanishes.

Further, in the thermodynamic limit the network cannot reach a disease-free steady state. In the thermodynamic limit, when $\gamma = 0$ a finite number of nodes will be infected. In this situation preferential attachment leads to a degree distribution with infinite variance, as node removal is so small. At the moment the variance of the degree distribution becomes infinite, we anticipate $\langle k_i \rangle$ to also become infinite and the limit p_c to go to zero. With the epidemic limit at zero the disease will be able to spread in the network.

We note that in finite networks the temporal dynamics can lead to a disease-free scale-free state. In this case the network goes through an initial exponential phase where an epidemic threshold exists that leads to the extinction of the epidemic. Subsequently, a scale-free topology develops, in which the epidemic threshold vanishes. However the epidemic cannot reappear as no infected are left which could reignite the epidemic. We emphasize that this mechanism is a finite-size effect. The finite-size extinction of the epidemic is demonstrated explicitly in Fig. 6, which shows that the fraction of individual-based simulation runs in which the epidemic persists increases with the initial network size M_0 . For populations that start with more than a few individuals the effect can be neglected and epidemics persist indefinitely.

Further dynamics. Precluding the finite-size effect described above, simulation runs in the parameter range considered so far approach a finite prevalence for any positive value of infectiousness. The paradox outlined above is thus resolved by alternative a) mentioned in the introduction, that is the model does not have an epidemic threshold although the second moment of the degree distribution remains finite. As a final step in our exploration present some evidence that the temporal solution of the paradox, alternative b), is also possible.

Up to now we have considered relatively small removal rates r . At sufficiently high removal rate, the evolution of the disease and the topology exhibits dynamics different from the convergence to a stationary state mainly due to the dominance of finite size effects. Figure 7 shows a representative evolution at non-zero fraction of infected arrivals w . At the combination of high infectiousness p and high removal rate r , the dynamics resembles a homoclinic trajectory in the $\langle k \rangle - [I]$ plane. The disease spreads quickly over the network and covers the whole population immediately. Then disease-induced removals dominate and the population becomes extinct until new healthy individuals arrive and the disease spreading restarts, which leads to cycles of population growth and collapse.

When all arrivals are susceptible, i.e. $w = 0$, the epidemics in the finite population disappears entirely and the collapse-and-growth cycle cannot be completed. A disease-free scale-free network then emerges. Based on the arguments above one can suspect that the same behavior cannot occur in the thermodynamic limit. Instead it is likely that the observed dynamics forms part of a homoclinic cycle, where long phases of very low disease prevalence are disrupted by sharp outbreaks.

The observed dynamics at high removal rates would correspond to diseases with very high mortality where infected individuals die almost immediately. This is reminiscent of the massive pandemics in history, where humanity was exposed to new pathogens with very high virulence. While such diseases could lead to the deaths of large fractions of populations, continuous supply of healthy individuals through immigration provided new hosts to the disease and introduced new bursts of disease spreading which caused repeated epidemic cycles.

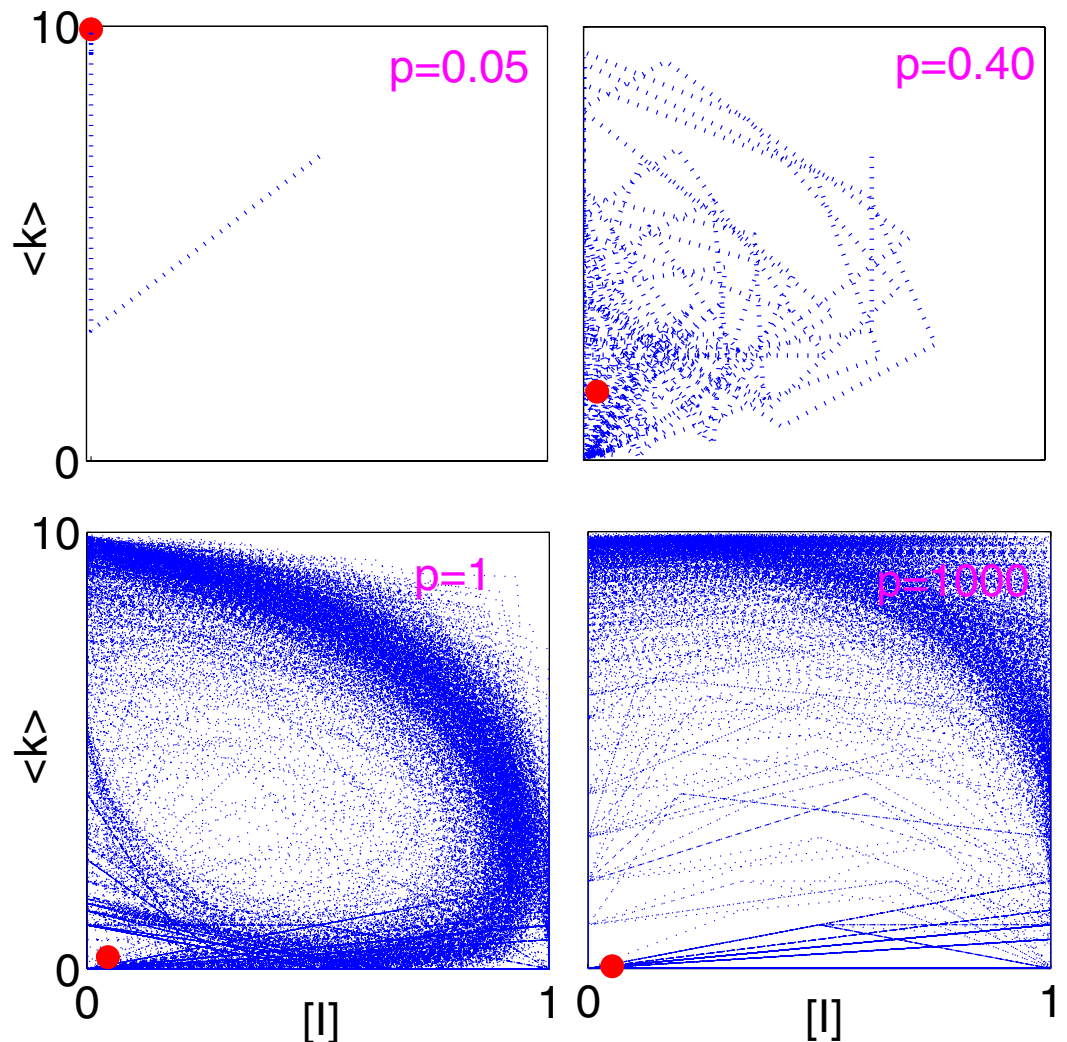


Figure 7. The dynamical evolution in the $[I] - \langle k \rangle$ plane for high removal rate r . At high infectiousness p , the observed behavior resembles a homoclinic trajectory. The network grows from a healthy initial state until an infected individual eventually arrives and the disease quickly spreads over the network and the infected individuals subsequently die. Homogeneous approximation (red circles) fails to capture this behavior and predicts a stable equilibrium. Parameters: $q = 0.01$, $r = 0.10$, $m_0 = 6$, $m = 5$, $w = 0.1$.

Discussion

In the present paper we have investigated the dynamics of a fatal SIR disease in a growing population. Our main finding is that no epidemic threshold exists in this model. Although the variance of the degree distribution remains finite in the evolved topologies topologies “unlikely” diseases with very low infectiousness can persist indefinitely.

We presented a detailed analytical exploration for the case of low removal rate, where the prevalence of the disease reaches a stationary level. In the growing population the ongoing epidemic dynamics eliminates the nodes of high degree and thus leads to the formation of topologies for which the variance of the degree distribution is finite. However, this mechanism only reduces the width of the degree distribution so far that the epidemic can still persist. Therefore, the disease itself can not reduce the degree variance so much to cause its own extinction. For any finite value of the infectiousness the network adapts its topology such that the variance of the degree distribution is lowered to point where the epidemic can still be sustained, which explains the observed absence of the epidemic threshold.

When the removal rate is sufficiently high, network simulations indicate that the disease can no longer stabilize at a stationary level and instead growth-and-collapse cycles are observed. Such dynamics are reminiscent of homoclinic orbits, which have for instance been observed in an adaptive-network model of social cooperation⁶⁸.

The dynamical feedback between the population structure and the epidemic disease has been so far studied in a number of articles in the past couple of years^{32,35–39,43,44,46,47}. However, models captured mostly social interactions in non-fatal diseases. A significant obstacle to progress in this line of work is that the network evolution in most models is driven by behavioral changes of individuals (whom to meet, how often to wash hands). However, despite the ever increasing availability of data, for behavior often no records exist such that model predictions

cannot easily be compared with real world data. By contrast, in the model proposed here, the social network evolves due to demographic processes such as migration and death on which data may be easier to obtain.

The model introduced here can be extended in several directions. For instance, dead hosts may not be immediately removed from the system and continue to spread the disease for a finite time until eventually being removed, the disease may not certainly lead to death such that infected individuals can recover, newly arriving individuals may establish links with only susceptible individuals, but not infected individuals, and susceptible individuals may rewire their links with infected neighbors to other susceptible individuals as a precaution against the epidemic. Independently of the specific mechanism, the disease-mortality combined with the network growth process can be expected to hold in all systems in which the dynamics of the population, in absence of the epidemic, leads to topologies with diverging variance of the degree distribution.

We believe that the proposed model is relevant for epidemics in rapidly growing cities, especially in the developing countries. In this context, connecting the model to real world data will be feasible in the future. In these situations we anticipate there is no epidemic threshold as the relevant populations often comprise many millions and are hence sufficiently large that the thermodynamic limit is a reasonable approximation, and the arrival of infected immigrants reintroduces diseases to the population. One can easily imagine extensions of the present model that incorporate policy measures such as vaccination, quarantine, or regulation of migration. We hope that this will in the future lead to the formulation of more efficient policies for combating epidemic diseases.

References

1. Hays, J. *Epidemics And Pandemics: Their Impacts on Human History* (ABC-CLIO, Santa Barbara, California, 2005).
2. WHO. *World Malaria Report* (WHO Press, World Health Organization, Geneva, Switzerland, 2011).
3. Organization, P. A. H. O. W. H. Zika suspected and confirmed cases reported by countries and territories in the americas cumulative cases. Tech. Rep., Pan American Health Organization 2015–2016.
4. Keeling, M. J. & Eames, K. T. D. Networks and epidemic models. *J. R. Soc. Interface* **2**, 295–307 (2005).
5. Bansal, S., Grenfell, B. T. & Meyers, L. A. When individual behaviour matters: homogeneous and network models in epidemiology. *J. R. Soc. Interface* **4**, 879–891 (2007).
6. Funk, S., Salathe, M. & Jansen, V. A. A. Modelling the influence of human behaviour on the spread of infectious diseases: a review. *J. R. Soc. Interface* **7**, 1247–1256 (2010).
7. Pastor-Satorras, R., Castellano, C., Mieghem, P. V. & Vespignani, A. Epidemic processes in complex networks. *Rev. Mod. Phys.* **87**, 925–979 (2015).
8. Castellano, C. & Pastor-Satorras, R. Thresholds for epidemic spreading in networks. *Phys. Rev. Lett.* **105**, 218701 (2010).
9. Boguñá, M., Castellano, C. & Pastor-Satorras, R. Nature of the epidemic threshold for the susceptible-infected-susceptible dynamics in networks. *Phys. Rev. Lett.* **111**, 068701 (2013).
10. Cohen, R., Erez, K., ben Avraham, D. & Havlin, S. Resilience of the internet to random breakdowns. *Phys. Rev. Lett.* **85**, 4626–4628 (2000).
11. Moreno, Y., Pastor-Satorras, R. & Vespignani, A. Epidemic outbreaks in complex heterogeneous networks. *Eur. Phys. J. B* **26**, 521–529 (2002).
12. Barthelemy, M., Barrat, A., Pastor-Satorras, R. & Vespignani, A. Velocity and hierarchical spread of epidemic outbreaks in scale-free networks. *Phys. Rev. Lett.* **92**, 178701 (2004).
13. Pastor-Satorras, R. & Vespignani, A. Immunization of complex networks. *Phys. Rev. E* **65**, 036104 (2002).
14. Zanette, D. H. & Kuperman, M. Effects of immunization in small-world epidemics. *Physica A* **309**, 445–452 (2002).
15. Cohen, R., Havlin, S. & Ben-avraham, D. Efficient immunization strategies for computer networks and populations. *Phys. Rev. Lett.* **91**, 247901 (2003).
16. Holme, P. Efficient local strategies for vaccination and network attack. *Europhys. Lett.* **68**, 908–914 (2004).
17. Meyers, L. A., Pourbohloul, B., Newman, M. E. J., Skowronski, D. M. & Brunham, R. C. Network theory and sars: predicting outbreak diversity. *Journal of Theoretical Biology* **232**, 71–81 (2005).
18. Hufnagel, L., Brockmann, D. & Geisel, T. Forecast and control of epidemics in a globalized world. *Proc. Natl. Acad. Sci. USA* **101**, 15124–15129 (2004).
19. Colizza, V., Barrat, A., Barthélemy, M. & Vespignani, A. The role of the airline transportation network in the prediction and predictability of global epidemics. *Proc. Natl. Acad. Sci. USA* **103**, 2015–2020 (2006).
20. Read, J. M., Eames, K. T. D. & Edmunds, W. J. Dynamic social networks and the implications for the spread of infectious disease. *J. R. Soc. Interface* **5**, 1001–1007 (2008).
21. Balcan, D. et al. Multiscale mobility networks and the spatial spreading of infectious diseases. *Proceedings of the National Academy of Sciences* **106**, 21484–21489 (2009).
22. Belik, V., Geisel, T. & Brockmann, D. Natural human mobility patterns and spatial spread of infectious diseases. *Phys. Rev. X* **1**, 011001 (2011).
23. Brockmann, D. & Helbing, D. The hidden geometry of complex, network-driven contagion phenomena. *Science* **342**, 1337–1342 (2013).
24. Pastor-Satorras, R. & Vespignani, A. Epidemic dynamics in finite size scale-free networks. *Phys. Rev. E* **65**, 035108(R) (2002).
25. Mieghem, P. V., Wang, H., Ge, X., Tang, S. & Kuipers, A. F. Influence of assortativity and degree-preserving rewiring on the spectra of networks. *The European Physical Journal B* **76**, 643–652 (2010).
26. Eguiluz, V. M. & Klemm, K. Epidemic threshold in structured scale-free networks. *Phys. Rev. Lett.* **89**, 108701 (2002).
27. Read, J. M. & Keeling, M. J. Disease evolution on networks: the role of contact structure. *Proc. R. Soc. B* **270**, 699–708 (2003).
28. Serrano, M. A. & Boguna, M. Percolation and epidemic thresholds in clustered networks. *Phys. Rev. Lett.* **97**, 088701 (2006).
29. Mossa, S., Barthélemy, M., Stanley, H. E. & Amaral, L. A. N. Truncation of power law behavior in scale-free network models due to information filtering. *Phys. Rev. Lett.* **88**, 138701 (2002).
30. Grabowski, A. & Kosinski, R. A. Epidemic spreading in a hierarchical social network. *Phys. Rev. E* **70**, 031908 (2004).
31. Newman, M. E. J. Spread of epidemic disease on networks. *Phys. Rev. E* **66**, 016128 (2002).
32. Gross, T., D’Lima, C. J. D. & Blasius, B. Epidemic dynamics on an adaptive network. *Phys. Rev. Lett.* **96**, 208701 (2006).
33. Gross, T. & Blasius, B. Adaptive coevolutionary networks: a review. *J. R. Soc. Interface* **5**, 259–271 (2008).
34. Gross, T. & Sayama, H. (eds) *Adaptive Networks. Theory, Models and Applications* (Springer Verlag, New York, USA, 2009).
35. Shaw, L. B. & Schwartz, I. B. Fluctuating epidemics on adaptive networks. *Phys. Rev. E* **77**, 066101 (2008).
36. Risau-Gusman, S. & Zanette, D. H. Contact switching as a control strategy for epidemic outbreaks. *Journal of Theoretical Biology* **257**, 52–60 (2009).
37. Wang, B., Cao, L., Suzuki, H. & Aihara, K. Epidemic spread in adaptive networks with multitype agents. *J. Phys. A* **44**, 035101 (2011).
38. Gräser, O., Hui, P. M. & Xu, C. Separatrices between healthy and endemic states in an adaptive epidemic model. *Physica A* **390**, 906–913 (2011).

39. Marceau, V., Noël, P.-A., Hebert-Dufresne, L., Allard, A. & Dube, L. Adaptive networks: coevolution of disease and topology. *Phys. Rev. E* **82**, 036116 (2010).
40. Shkarayev, M. S., Tunc, I. & Shaw, L. B. Epidemics with temporary link deactivation in scale-free networks. *Journal of Physics A: Mathematical and Theoretical* **47**, 455006 (2014).
41. Shaw, L. B. & Schwartz, I. B. Enhanced vaccine control of epidemics in adaptive networks. *Phys. Rev. E* **81**, 046120 (2010).
42. Lagorio, C. *et al.* Quarantine-generated phase transition in epidemic spreading. *Phys. Rev. E* **83**, 026102 (2011).
43. Zanette, D. H. & Risau-Gusmán, S. Infection Spreading in a Population with Evolving Contacts. *Journal of Biological Physics* **34**, 135–148 URL <http://dx.doi.org/10.1007/s10867-008-9060-9> (2008).
44. Yang, H., Tang, M. & Zhang, H.-F. Efficient community-based control strategies in adaptive networks. *New Journal of Physics* **14**, 123017 (2012).
45. Tunc, I., Shkarayev, M. S. & Shaw, L. B. Epidemics in adaptive social networks with temporary link deactivation. *Journal of Statistical Physics* **151**, 355–366 (2013).
46. Zhang, H.-F., Xie, J.-R., Tang, M. & Lai, Y.-C. Suppression of epidemic spreading in complex networks by local information based behavioral responses. *Chaos* **24** (2014).
47. Yang, H., Tang, M. & Gross, T. Large epidemic thresholds emerge in heterogeneous networks of heterogeneous nodes. *Scientific Reports* **5**, 13122 (2015).
48. Scarpino, S. V., Allard, A. & Hebert-Dufresne, L. The effect of a prudent adaptive behaviour on disease transmission. *Nat Phys* advance online publication (2016).
49. Kamenov, A. & Meerson, B. Extinction of an infectious disease: a large fluctuation in a nonequilibrium system. *Phys. Rev. E* **77**, 061107 (2008).
50. Schwartz, I. B., Billings, L., Dykman, M. & Landsman, A. Predicting extinction rates in stochastic epidemic models. *J. Stat. Mech.* P01005 (2009).
51. Guerra, B. & Gómez-Gardeñes, J. Annealed and mean-field formulations of disease dynamics on static and adaptive networks. *Phys. Rev. E* **82**, 035101(R) (2010).
52. Althouse, B. M. & Hébert-Dufresne, L. Epidemic cycles driven by host behaviour. *Journal of The Royal Society Interface* **11** (2014).
53. Hayashi, Y., Minoura, M. & Matsukubo, J. Oscillatory epidemic prevalence in growing scale-free networks. *Phys. Rev. E* **69**, 016112 (2004).
54. Poncela, J., Gómez-Gardeñes, J., Traulsen, A. & Moreno, Y. Evolutionary game dynamics in a growing structured population. *New J. Phys.* **11**, 083031 (2009).
55. Anderson, R. M. & May, R. M. (eds) *Infectious Diseases of Humans: Dynamics and Control* (Oxford University Press, New York, USA, 1991).
56. Barabási, A.-L. & Albert, R. Emergence of scaling in random networks. *Science* **286**, 509–512 (1999).
57. Keeling, M. J. The effects of local spatial structure on epidemiological invasions. *Proc. R. Soc. B* **266**, 859–867 (1999).
58. Bauch, C. T. The spread of infectious diseases in spatially structured populations: an invasy pair approximation. *Mathematical Biosciences* **198**, 217–237 (2005).
59. Noël, P.-A., Davoudi, B., Brunham, R. C., Dube, L. J. & Pourbohloul, B. Time evolution of epidemic disease on finite and infinite networks. *Phys. Rev. E* **79**, 026101 (2009).
60. House, T. & Keeling, M. J. Insights from unifying modern approximations to infections on networks. *J. R. Soc. Interface* **8**, 67–73 (2011).
61. Lindquist, J., Ma, J., van den Driessche, P. & Willeboordse, F. H. Effective degree network disease models. *J. Math. Biol.* **62**, 143–164 (2011).
62. Gleeson, J. P. High-accuracy approximation of binary-state dynamics on networks. *Phys. Rev. Lett.* **107**, 068701 (2011).
63. Demirel, G., Vazquez, F., Böhme, G. A. & Gross, T. Moment-closure approximations for discrete adaptive networks. *Physica D* (2013).
64. Gleeson, J. P. Binary-state dynamics on complex networks: Pair approximation and beyond. *Phys. Rev. X* **3**, 021004 (2013).
65. Pastor-Satorras, R. & Vespignani, A. Epidemic spreading in scale-free networks. *Phys. Rev. Lett.* **86**, 3200–3203 (2001).
66. Pastor-Satorras, R. & Vespignani, A. Epidemic dynamics and endemic states in complex networks. *Phys. Rev. E* **63**, 066117 (2001).
67. Silk, H., M. H., Demirel, G. & Gross, T. Exploring the adaptive voter model dynamics with a mathematical triple jump. *New Journal of Physics* **16**, 93051 (2014).
68. Zschaler, G., Traulsen, A. & Gross, T. A homoclinic route to asymptotic full cooperation in adaptive networks and its failure. *New J. Phys.* **12**, 093015 (2010).

Acknowledgements

This work was supported partially by the EPSRC under grant code EP/I013717/1(BCCS) and a Max-Planck Society PhD scholarship.

Author Contributions

All authors contributed to and reviewed the manuscript. G.D prepared the figures.

Additional Information

Competing financial interests: The authors declare no competing financial interests.

How to cite this article: Demirel, G. *et al.* Dynamics of epidemic diseases on a growing adaptive network. *Sci. Rep.* **7**, 42352; doi: 10.1038/srep42352 (2017).

Publisher's note: Springer Nature remains neutral with regard to jurisdictional claims in published maps and institutional affiliations.



This work is licensed under a Creative Commons Attribution 4.0 International License. The images or other third party material in this article are included in the article's Creative Commons license, unless indicated otherwise in the credit line; if the material is not included under the Creative Commons license, users will need to obtain permission from the license holder to reproduce the material. To view a copy of this license, visit <http://creativecommons.org/licenses/by/4.0/>

© The Author(s) 2017